

Transparents du huitième cours

- **Modèle Standard**
- γ_5 et **Dim-Reg**
- **Espaces de dimension z**
- **Anomalies et formule de l'indice locale**

Modèle Standard

$$\begin{aligned}
\mathcal{L}_{SM} = & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \frac{1}{2}ig_s^2 (\bar{q}_i^\sigma \gamma^\mu q_j^\sigma) g_\mu^a \\
& + \bar{G}^a \partial^2 G^a + g_s f^{abc} \partial_\mu \bar{G}^a G^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 \\
& - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- \\
& - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h \left(\frac{2M^2}{g^2} + \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right) \\
& + \frac{2M^4}{g^2} \alpha_h - igc_w (\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) \\
& \quad + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)) - igs_w (\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) \\
& \quad - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)) \\
& - \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\mu^0 W_\nu^+ W_\nu^-) \\
& + g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\mu W_\nu^+ W_\nu^-) + g^2 s_w c_w (A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) \\
& \quad - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-) - g\alpha_h M (H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-) \\
& - \frac{1}{8}g^2 \alpha_h (H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2) \\
& \quad - gMW_\mu^+ W_\mu^- H - \frac{1}{2}g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H \\
& - \frac{1}{2}ig (W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0))
\end{aligned}$$

$$\begin{aligned}
& +\frac{1}{2}g \left(W_\mu^+ (H\partial_\mu\phi^- - \phi^-\partial_\mu H) - W_\mu^- (H\partial_\mu\phi^+ - \phi^+\partial_\mu H) \right) \\
& +\frac{1}{2}g\frac{1}{c_w} (Z_\mu^0 (H\partial_\mu\phi^0 - \phi^0\partial_\mu H) - ig\frac{s_w^2}{c_w} MZ_\mu^0 (W_\mu^+\phi^- - W_\mu^-\phi^+)) \\
& +igs_w MA_\mu (W_\mu^+\phi^- - W_\mu^-\phi^+) - ig\frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) \\
& +igs_w A_\mu (\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- (H^2 + (\phi^0)^2 + 2\phi^+\phi^-) \\
& \quad - \frac{1}{4}g^2 \frac{1}{c_w^2} Z_\mu^0 Z_\mu^0 (H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2\phi^+\phi^-) \\
& -\frac{1}{2}g^2 \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+\phi^- + W_\mu^-\phi^+) - \frac{1}{2}ig^2 \frac{s_w^2}{c_w} Z_\mu^0 H (W_\mu^+\phi^- - W_\mu^-\phi^+) \\
& +\frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+\phi^- + W_\mu^-\phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+\phi^- - W_\mu^-\phi^+) \\
& \quad -g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+\phi^- - g^2 s_w^2 A_\mu A_\mu \phi^+\phi^- \\
& -\bar{e}^\lambda (\gamma\partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda \gamma\partial\nu^\lambda - \bar{u}_j^\lambda (\gamma\partial + m_u^\lambda) u_j^\lambda - \bar{d}_j^\lambda (\gamma\partial + m_d^\lambda) d_j^\lambda \\
& \quad +igs_w A_\mu \left(-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda) \right) \\
& \quad +\frac{ig}{4c_w} Z_\mu^0 \{ (\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) \\
& \quad +(\bar{u}_j^\lambda \gamma^\mu (\frac{4}{3}s_w^2 - 1 - \gamma^5) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu (1 - \frac{8}{3}s_w^2 - \gamma^5) d_j^\lambda) \} \\
& \quad +\frac{ig}{2\sqrt{2}} W_\mu^+ \left((\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\kappa} d_j^\kappa) \right) \\
& \quad +\frac{ig}{2\sqrt{2}} W_\mu^- \left((\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger \gamma^\mu (1 + \gamma^5) u_j^\lambda) \right)
\end{aligned}$$

$$\begin{aligned}
& + \frac{ig}{2\sqrt{2}} \frac{m_e^\lambda}{M} \left(-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda) \right) \\
& \quad - \frac{g}{2} \frac{m_e^\lambda}{M} \left(H(\bar{e}^\lambda e^\lambda) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda) \right) \\
& + \frac{ig}{2M\sqrt{2}} \phi^+ \left(-m_d^\kappa (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + m_u^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\kappa) \right) \\
& + \frac{ig}{2M\sqrt{2}} \phi^- \left(m_d^\lambda (\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger (1 + \gamma^5) u_j^\kappa) - m_u^\kappa (\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger (1 - \gamma^5) u_j^\kappa) \right) \\
& - \frac{g}{2} \frac{m_u^\lambda}{M} H(\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2} \frac{m_d^\lambda}{M} H(\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2} \frac{m_u^\lambda}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \frac{ig}{2} \frac{m_d^\lambda}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) \\
& + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \frac{M^2}{c_w^2}) X^0 \\
& + \bar{Y} \partial^2 Y + igc_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + igs_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ Y) \\
& + igc_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + igs_w W_\mu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) \\
& + igc_w Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + igs_w A_\mu (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) \\
& \quad - \frac{1}{2} gM \left(\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w^2} \bar{X}^0 X^0 H \right) \\
& \quad + \frac{1 - 2c_w^2}{2c_w} igM \left(\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^- \right) \\
& \quad + \frac{1}{2c_w} igM \left(\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^- \right) \\
& \quad + igMs_w \left(\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^- \right) \\
& \quad + \frac{1}{2} igM \left(\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0 \right) .
\end{aligned}$$

Notations

Bosons de jauge : $A_\mu, W_\mu^\pm, Z_\mu^0, g_\mu^a$

Quarks : u_j^κ, d_j^κ et $q_j^\sigma =$ collectif pour quarks

Leptons : e^λ, ν^λ

Higgs : $H, \phi^0, \phi^+, \phi^-$

Fantômes : G^a, X^0, X^+, X^-, Y , jauge Feynman

Masses : $m_d^\lambda, m_u^\lambda, m_e^\lambda, m_h, M \leftarrow$ Masse du W

Constantes de couplage : $g = \sqrt{4\pi\alpha}$ (structure fine), $g_s =$ forte, $\alpha_h = \frac{m_h^2}{4M^2}$

Constante de tadpole β_h

Sinus et cosinus de l'angle faible : s_w, c_w

Cabibbo-Kobayashi-Maskawa : $C_{\lambda\kappa}$

Constantes de structure de SU_3 : f^{abc}

Indice Local NCG (ac+hm)

$$\int P := \text{Res}_{z=0} \text{Tr} (P |D|^{-z})$$

Trace sur l'algèbre engendrée par \mathcal{A} , $[D, \mathcal{A}]$ et $|D|^z$, $z \in \mathbb{C}$.

$$\varphi_0(a) = \lim_{z \rightarrow 0} \text{Tr}(\gamma a |D|^{-z}), \quad \forall a \in \mathcal{A},$$

$$\varphi_n(a^0, \dots, a^n) :=$$

$$\sum_k c_{n,k} \int \gamma a^0 [D, a^1]^{(k_1)} \dots [D, a^n]^{(k_n)} |D|^{-n-2|k|}$$

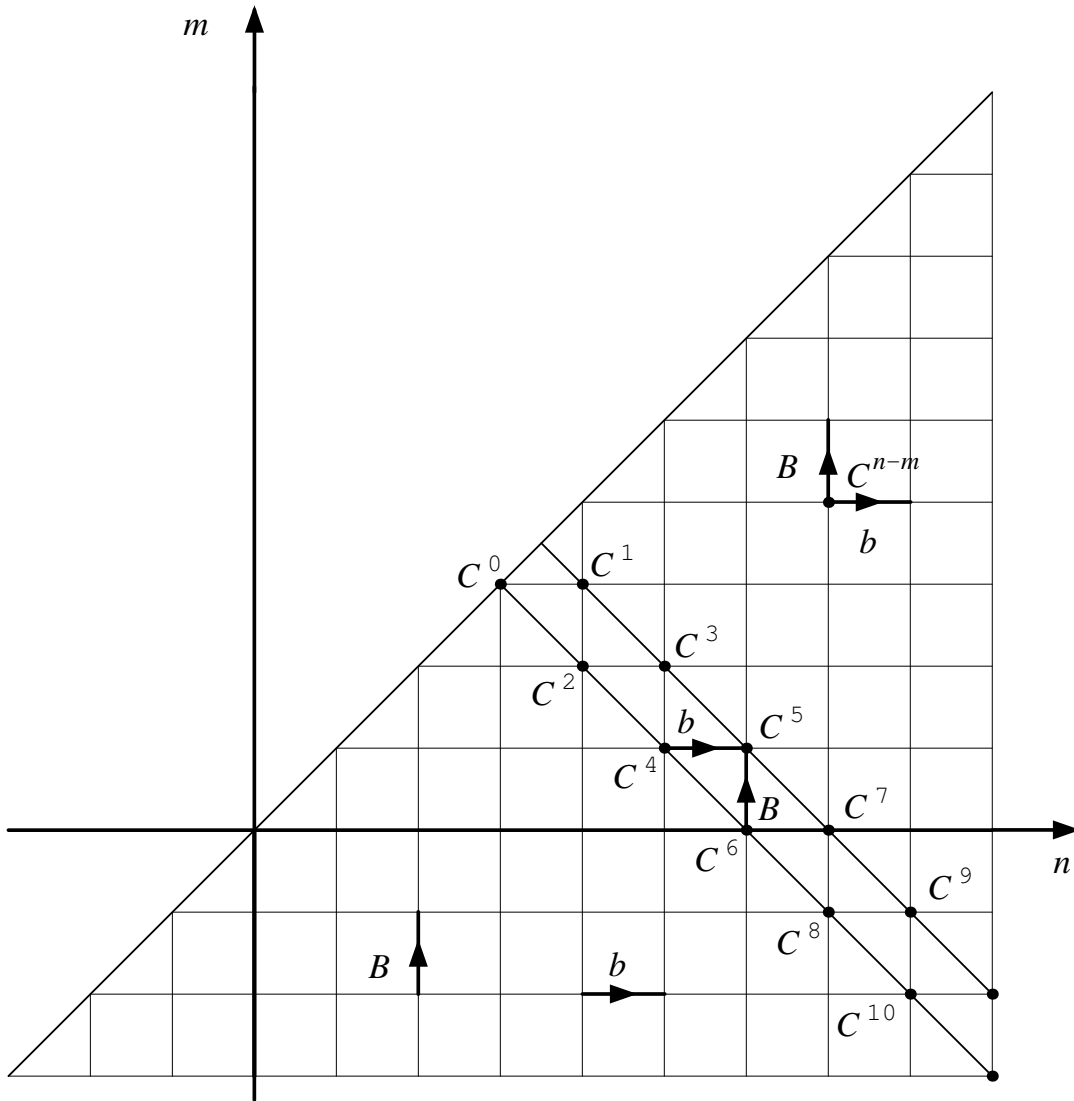
$T^{(k_i)} = \nabla^{k_i}(T)$ avec $\nabla(T) = D^2T - TD^2$, $|k| = k_1 + \dots + k_n$,

$$c_{n,k} = \frac{(-1)^{|k|}}{2} (k_1! \dots k_n!)^{-1}$$

$$((k_1+1) \dots (k_1+k_2+\dots+k_n+n))^{-1} \Gamma(|k| + n/2).$$

$(\varphi_n)_{n=0,2,\dots}$ cocycle du (b, B) -bicomplexe de \mathcal{A} .

Accouplement $(\varphi_n) \in HC^*(\mathcal{A})$ avec $K_0(\mathcal{A}) =$ indice D , $K_0(\mathcal{A}) \rightarrow \mathbb{Z}$.



Dim-Reg

Espaces X_z de dimension z (ac + mm)

t'Hooft-Veltman et Breitenlohner-Maison \Leftrightarrow faire le produit de l'espace-temps euclidien par un triplet spectral X_z de dimension $z \in \mathbb{C}$, $\text{Re}(z) > 0$

$$\mathcal{H}'' = \mathcal{H} \otimes \mathcal{H}', \quad D'' = D \otimes 1 + \gamma_5 \otimes D_z$$

Spectre de dimensions de X_z est réduit à z .

$$\text{Trace}(e^{-\lambda D_z^2}) = \pi^{z/2} \lambda^{-z/2}, \quad \forall \lambda \in \mathbb{R}_+^*$$

Espaces X_z

$$\mathrm{Tr}_N(1_E(Z)) = \frac{1}{2} \int_E dy$$

$$D_z = \rho(z) F |Z|^{1/z}$$

$$\rho(z) = \pi^{-\frac{1}{2}} \left(\Gamma\left(\frac{z}{2} + 1\right) \right)^{\frac{1}{z}}$$

L'opérateur D_z vérifie pour $z > 0$

$$\mathrm{Tr}_N(e^{-\lambda D_z^2}) = \pi^{z/2} \lambda^{-z/2}, \quad \forall \lambda \in \mathbb{R}_+^*$$

et pour $z \neq 0$,

$$\mathrm{Tr}_N((D_z^2)^{s/2})$$

a un pôle simple en $s = z$ et est absolument convergente dans $\mathrm{Re}(s/z) > 1$.

Cutoff infrarouge

$$\begin{aligned}\mathrm{Tr}'_N((D_z^2)^{-s/2}) &= \frac{1}{2} \int_{|y|>1} (\rho^2 |y|^{2/z})^{-s/2} dy \\ &= \rho^{-s} \int_1^\infty u^{-s/z} du = \rho^{-s} \frac{z}{s-z}\end{aligned}$$

$$R(\lambda, z) = \int_0^{\frac{1}{2}} (e^{-\lambda \rho^2 |y|^{2/z}} - e^{-\lambda \rho^2 f(|y|)^{2/z}}) dy$$

$$|R(\lambda, z)| < C |\lambda| 2^{\mathrm{Re}(-2/z)}$$

Potentiel de jauge

$(\mathcal{A}, \mathcal{H}, D)$ et \mathcal{E} projectif de type fini sur \mathcal{A} ,

$\mathcal{B} = \text{End}_{\mathcal{A}}(\mathcal{E})$, $\mathcal{H}' = \mathcal{E} \otimes_{\mathcal{A}} \mathcal{H}$, $D' = ?$

$$D'(\xi \otimes \eta) = \xi \otimes D\eta + \nabla(\xi)\eta$$

$$\nabla(\xi a) = (\nabla\xi)a + \xi \otimes da, \quad da = [D, a]$$

$$A = \sum a_i [D, b_i], \quad a_i, b_i \in \mathcal{A} \rightarrow \Omega_D^1 \subset \mathcal{L}(\mathcal{H})$$

$$D \mapsto D + A, \quad A = A^*$$

$$u(D + A)u^* = D + \alpha_u(A),$$

$$\alpha_u(A) = u[D, u^*] + u A u^*$$

Potentiel évanescent

\mathcal{A} algèbre $\mathbb{Z}/2$ -graduée,

$$[D, a]_- := D a - (-1)^{\deg(a)} a D$$

$$\tilde{\mathcal{A}} = \{a + b\gamma\} \sim \mathcal{A} \oplus \mathcal{A}, a + b\gamma \mapsto (a + b, a - b)$$

$\mathbb{Z}/2$ -graduée par

$$\theta \in \text{Aut}(\tilde{\mathcal{A}}), \quad \theta(\gamma) = -\gamma, \quad \theta(a) = a, \quad \forall a \in \mathcal{A}$$

$$\bar{D} = D \otimes 1, \quad \hat{D} = \gamma \otimes D_z, \quad D'' = \bar{D} + \hat{D}$$

Potentiel évanescent

$$B = [D'', \gamma]_- = 2\gamma \hat{D}$$

$$E = \frac{1}{2} a [D'', \gamma]_- = \gamma a \hat{D}$$

Intégrale fonctionnelle fermionique

Euclidien

$$\mathcal{L}_{\text{fermions}} = \langle \eta, D'' \xi \rangle$$

avec deux variables *indépendantes* ξ et η .

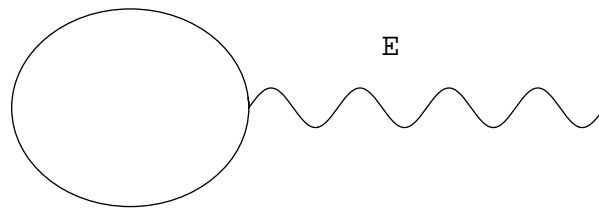
Transformations de jauge graduées

$$\alpha_u(\xi) = u \xi, \quad \alpha_u(\eta) = \theta(u) \eta$$

Symétrie chirale $u = e^{ia\gamma}$

$$\delta \mathcal{L}_{\text{fermions}} = \langle \eta, (i [D, a] \gamma + 2 i a \gamma \hat{D}) \xi \rangle$$

Tadpole



$$E = \gamma a \hat{D}$$

Soit φ_0 la composante de dimension 0 du cocycle local, alors (quand $z \rightarrow 0$)

$$\text{Tr}(E D''^{-1}) = -\varphi_0(a)$$

$$D''^{-1} = D'' D''^{-2} = (\bar{D} + \hat{D}) D''^{-2}$$

$$D''^2 = \bar{D}^2 + \hat{D}^2$$

$$D''^{-2} = \int_0^\infty e^{-t\bar{D}^2} e^{-t\hat{D}^2} dt$$

$$\text{Tr}(E D''^{-1}) = \text{Tr}(\gamma a \hat{D} (\bar{D} + \hat{D}) D''^{-2}) =$$

$$\text{Tr}(\gamma a \hat{D}^2 D''^{-2}) =$$

$$\int_0^\infty \text{Tr}(\gamma a \hat{D}^2 e^{-t\bar{D}^2} e^{-t\hat{D}^2}) dt =$$

$$\int_0^\infty \text{Tr}(\gamma a e^{-tD^2}) \text{Tr}(D_z^2 e^{-tD_z^2}) dt$$

$$\text{Tr}_N(D_z^2 e^{-tD_z^2}) = \frac{z}{2} \pi^{z/2} t^{-z/2-1} \quad \forall t \in \mathbb{R}_+^*$$

$$\int_0^\infty e^{-tD^2} t^{-z/2-1} dt = \Gamma(-z/2) |D|^z$$

Opérateurs différentiels : $OP(\mathcal{A}, \mathcal{H}, D)$

$$D^{-2} T \sim \sum_0^{\infty} (-1)^k \nabla^k(T) D^{-2k-2}$$

Lemme

Soit $P \in OP(\mathcal{A}, \mathcal{H}, D)$. Pour $n > k > 0$ et $z \rightarrow 0$,

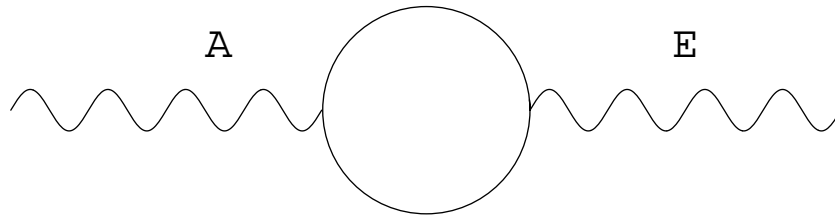
$$\begin{aligned} \text{Tr} (\gamma \hat{D}^{2k} (P \otimes 1) D''^{-2n}) = \\ -\frac{1}{2} B(k, n - k) \int \gamma P D^{-2(n-k)} \end{aligned}$$

$$\begin{aligned}
& \text{Tr}(\gamma \hat{D}^{2k} (P \otimes 1) D''^{-2n}) = \\
& \frac{1}{\Gamma(n)} \int_0^\infty \text{Tr}(\gamma \hat{D}^{2k} (P \otimes 1) e^{-t\bar{D}^2} e^{-t\hat{D}^2}) t^{n-1} dt \\
& \text{Tr}_N(D_z^{2k} e^{-tD_z^2}) = \frac{\prod_0^{k-1} (z + 2j)}{2^k} \pi^{z/2} t^{-z/2-k} \\
& \int_0^\infty e^{-tD^2} t^{n-1-z/2-k} dt = \Gamma(n-z/2-k) |D|^{z-2(n-k)} \\
& \oint Q := \text{Res}_{z=0} \text{Tr}(Q|D|^{-z})
\end{aligned}$$

$$Q = \gamma P D^{-2(n-k)}$$

$$B(p, q) = \frac{\Gamma(p) \Gamma(q)}{\Gamma(p+q)}$$

Self-énergie



$$\text{Tr}(E D''^{-1} A D''^{-1}) =$$

$$\sum_0^{\infty} (-1)^{n+1} \frac{1}{2n+2} \int \gamma a \nabla^n(B) D^{-2n-2}$$

où $B = DA + AD = dA + A'$ avec

$$dA = \sum [D, a_i] [D, b_i], \quad A' = \sum a_i \nabla(b_i)$$

$$\begin{aligned}
& \text{Tr}(E D''^{-1} A D''^{-1}) = \\
& \text{Tr}(\gamma a \hat{D} (\bar{D} + \hat{D}) D''^{-2} A (\bar{D} + \hat{D}) D''^{-2}) = \\
& \quad \text{Tr}(\gamma a \hat{D}^2 D''^{-2} A \bar{D} D''^{-2}) + \\
& \quad \text{Tr}(\gamma a \hat{D} \bar{D} D''^{-2} A \hat{D} D''^{-2}) = \\
& \quad \text{Trace}(\gamma a \hat{D}^2 D''^{-2} (A \bar{D} + \bar{D} A) D''^{-2}) \\
& (A \bar{D} + \bar{D} A) = B \otimes 1, \quad B = dA + A' \\
& D^{-2} T \sim \sum_0^{\infty} (-1)^k \nabla^k(T) D^{-2k-2} \\
& \text{Trace}(E D''^{-1} A D''^{-1}) = \\
& \sum_0^{\infty} (-1)^n \text{Tr}(\gamma a \hat{D}^2 (\nabla^n(B) \otimes 1) D''^{-2n-4}).
\end{aligned}$$

$$D^2 T D^{-2} = T + \nabla(T) D^{-2} = (1 + \epsilon)(T)$$

$$\Theta(T) = \sum_1^{\infty} \frac{(-1)^{n+1}}{n} \nabla^n(T) D^{-2n}$$

$$\pi(z) = \frac{z}{e^z - 1}, \quad \pi(-z) = e^z \pi(z)$$

$$\text{Tr}(E D''^{-1} A D''^{-1}) = -\frac{1}{2} \int \gamma a \pi(\Theta)(B) D^{-2} =$$

$$-\frac{1}{2} \int \gamma A \pi(\Theta)([D, a]) D^{-2}$$

$$\int X \pi(\Theta)(Y) D^{-2} = \int Y \pi(\Theta)(X) D^{-2}$$

Self-énergie en $D = 2$

Cocycle de Hochschild $\varphi \rightarrow \int_{\varphi}$

$$\int_{\varphi} a_0 da_1 \cdots da_n = \varphi(a_0, a_1, \cdots, a_n)$$

$$\int_{\varphi} a \omega = \int_{\varphi} \omega a, \quad \forall a \in \mathcal{A}$$

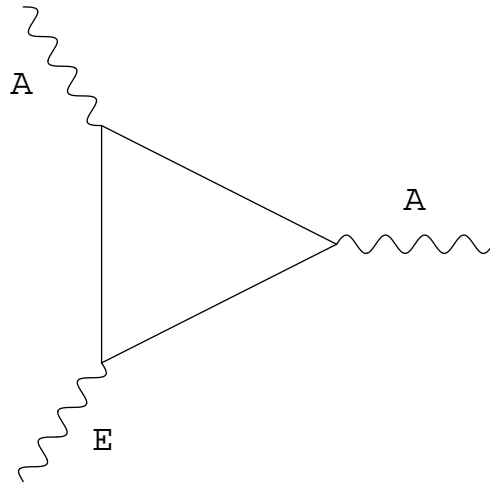
$$\varphi_2(a_0, a_1, a_2) = \frac{1}{4} \int \gamma a_0 [D, a_1] [D, a_2] D^{-2}$$

$$\text{Tr}(E D''^{-1} A D''^{-1}) = -2 \int_{\varphi_2} A da$$

Tadpole = 0 $\Rightarrow \varphi_2$ cyclique \Rightarrow

$$\text{Tr}(E D''^{-1} A D''^{-1}) = -2 \int_{\varphi_2} a dA$$

ABJ triangle



$$\text{Tr}(E D''^{-1} A D''^{-1} A D''^{-1})$$

ABJ en $D = 2$

$$\text{Tr}(E D''^{-1} A D''^{-1} A D''^{-1}) = 2 \int_{\varphi_2} a A^2$$

1. φ_2 est un deux cocycle cyclique* dans la classe locale.
2. Pour tout $a \in \mathcal{A}$ et A

$$\begin{aligned} \text{Tr}(E D''^{-1} A D''^{-1} A D''^{-1}) - \text{Tr}(E D''^{-1} A D''^{-1}) \\ = 2 \int_{\varphi_2} a (dA + A^2) \end{aligned}$$

*on suppose tadpole= 0

$$\begin{aligned} & \text{Tr}(E D''^{-1} A D''^{-1} A D''^{-1}) = \\ & \text{Tr}(\gamma a \hat{D} (\bar{D} + \hat{D}) D''^{-2} A (\bar{D} + \hat{D}) D''^{-2} A (\bar{D} + \hat{D}) D''^{-2}) \end{aligned}$$

Nombre impair de \hat{D} donne 0

Termes avec 4 fois \hat{D} ,

$$\begin{aligned} T_4 = & \text{Tr}(\gamma a \hat{D}^2 D''^{-2} A \hat{D} D''^{-2} A \hat{D} D''^{-2}) = \\ & -\text{Tr}(\gamma a \hat{D}^4 D''^{-2} A D''^{-2} A D''^{-2}) \end{aligned}$$

car $\hat{D} = \gamma \otimes D_z$ anticommute avec A (à cause de γ)

Termes en \hat{D}^2

$$T_1 = \text{Tr}(\gamma_a \hat{D}^2 D''^{-2} A \bar{D} D''^{-2} A \bar{D} D''^{-2})$$

$$\begin{aligned} T_2 &= \text{Tr}(\gamma_a \hat{D} \bar{D} D''^{-2} A \hat{D} D''^{-2} A \bar{D} D''^{-2}) = \\ &\text{Tr}(\gamma_a \hat{D}^2 \bar{D} D''^{-2} A D''^{-2} A \bar{D} D''^{-2}) \end{aligned}$$

$$\begin{aligned} T_3 &= \text{Tr}(\gamma_a \hat{D} \bar{D} D''^{-2} A \bar{D} D''^{-2} A \hat{D} D''^{-2}) = \\ &\text{Tr}(\gamma_a \hat{D}^2 \bar{D} D''^{-2} A \bar{D} D''^{-2} A D''^{-2}) \end{aligned}$$

Lemme

$$\text{Tr}(\hat{D}^{2k} P_0 D''^{-2} P_1 D''^{-2} P_2 D''^{-2}) =$$

$$\sum c(a, b, k) \int P_0 \nabla^a(P_1) \nabla^b(P_2) D''^{-2(a+b+3-k)}$$

$$c(a, b, k) = (-1)^{a+b+1} \frac{1}{2} \frac{((k-1)!(a+b+2-k)!)}{b!(a+1)!(a+b+2)}$$

$$D''^{-2} P_1 D''^{-2} P_2 \sim \sum d(a, b) \nabla^a(P_1) \nabla^b(P_2) D''^{-2(a+b+2)}$$

avec

$$d(a, b) = (-1)^{a+b} \sum_{0 \leq c \leq b} \frac{(a+c)!}{a! c!} =$$

$$(-1)^{a+b} \frac{(a+b+1)!}{b! (a+1)!}$$

$$\frac{1}{2} \frac{(a+b+1)! (k-1)! (a+b+2-k)!}{b! (a+1)! (a+b+2)!} =$$

$$\frac{1}{2} \frac{((k-1)! (a+b+2-k)!)}{b! (a+1)! (a+b+2)}$$

$$\begin{aligned}
T_4 &= -\text{Tr}(\gamma_a \hat{D}^4 D''^{-2} A D''^{-2} A D''^{-2}) \\
&= \frac{1}{4} \int \gamma_a A^2 D^{-2}
\end{aligned}$$

$$\begin{aligned}
T_2 &= \text{Tr}(\gamma_a \hat{D}^2 \bar{D} D''^{-2} A D''^{-2} A \bar{D} D''^{-2}) \\
&= -\frac{1}{4} \int \gamma_a D A^2 D^{-3} = \frac{1}{4} \int \gamma_a A^2 D^{-2}
\end{aligned}$$

↓

$$T_2 + T_4 = 2 \int_{\varphi_2} a A^2$$

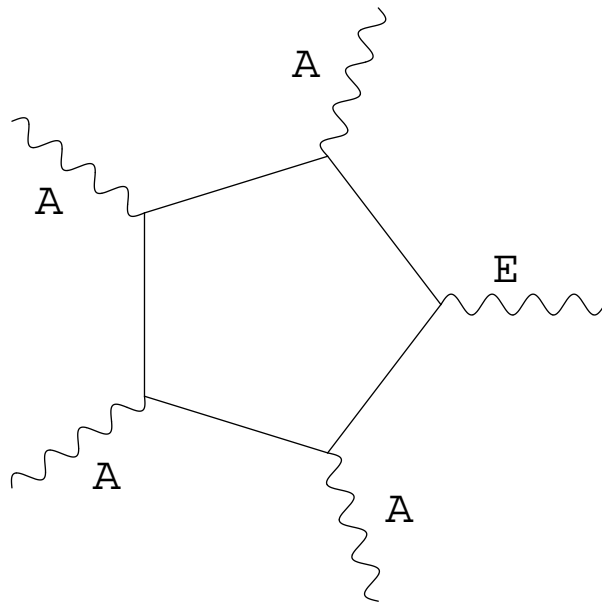
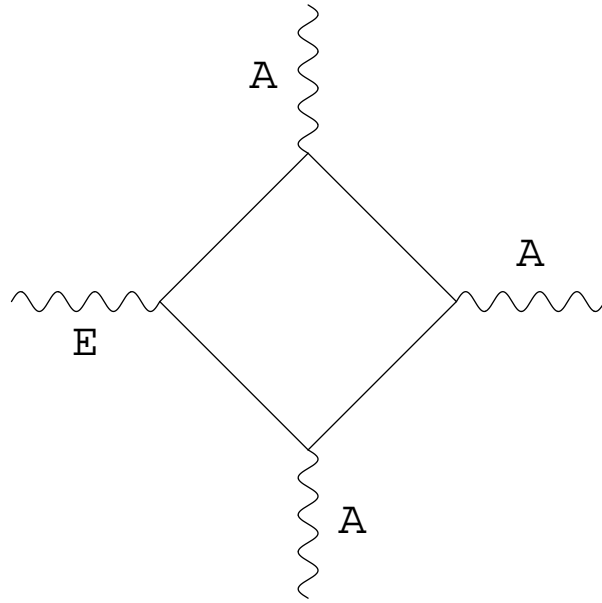
$$\begin{aligned}
T_1 &= \text{Tr}(\gamma_a \hat{D}^2 D''^{-2} A \bar{D} D''^{-2} A \bar{D} D''^{-2}) \\
&= -\frac{1}{4} \int \gamma_a A D A D^{-3}
\end{aligned}$$

$$\begin{aligned}
T_3 &= \text{Tr}(\gamma_a \hat{D}^2 \bar{D} D''^{-2} A \bar{D} D''^{-2} A D''^{-2}) \\
-\frac{1}{4} \int \gamma_a D A D A D^{-4} &= -\frac{1}{4} \int \gamma D a A D A D^{-4} \\
&= \frac{1}{4} \int \gamma_a A D A D^{-3} = -T_1
\end{aligned}$$

↓

$$T_1 + T_3 = 0$$

$$D = 4$$



Fermions \rightarrow Géométrie

Fermions	$\psi \in \mathcal{H}$
Symétries internes	Int(\mathcal{A}) $f \rightarrow u f u^*$
Bosons de Jauge	Fluctuations internes

Action Bosonique = Action Spectrale
(ac+ac)

$N(\Lambda) = \#$ valeurs propres de D dans $[-\Lambda, \Lambda]$.

$$N(\Lambda) = \langle N(\Lambda) \rangle + N_{\text{osc}}(\Lambda)$$

$$\langle N(\Lambda) \rangle = S_{\Lambda}(D) = \sum_{k \in S} \frac{\Lambda^k}{k} \int |ds|^k + \zeta_D(0),$$

$$\zeta_D(s) = \text{Trace}(|D|^{-s})$$

Modèle Standard en couplage minimal

$$\mathcal{L}_E + \mathcal{L}_G + \mathcal{L}_{GH} + \mathcal{L}_H + \mathcal{L}_{Gf} + \mathcal{L}_{Hf}$$

Action Spectrale (ac+ac)

$$\begin{aligned} S = & \int d^4x \sqrt{g} \left(\frac{1}{2\kappa_0^2} R - \mu_0^2 (H^* H) \right. \\ & + a_0 C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} + b_0 R^2 + c_0 {}^*R^*R + d_0 R_{;\mu}{}^\mu \\ & + e_0 + \frac{1}{4} G_{\mu\nu}^i G^{\mu\nu i} + \frac{1}{4} F_{\mu\nu}^\alpha F^{\mu\nu\alpha} \\ & \left. + \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + |D_\mu H|^2 - \xi_0 R |H|^2 + \lambda_0 (H^* H)^2 \right) \end{aligned}$$

Modèle Standard

$$X = M \times F$$

$$\mathcal{A} = \mathcal{A}_M \otimes \mathcal{A}_F, \quad \mathcal{H} = \mathcal{H}_M \otimes \mathcal{H}_F,$$

$$D = D_M \otimes 1 + \gamma_5 \otimes D_F$$

$$\mathcal{A}_F = \mathbb{C} \oplus \mathbb{H} \oplus M_3(\mathbb{C})$$

$$\mathcal{H}_F = Q \oplus L \oplus \bar{Q} \oplus \bar{L}$$

$$Q = \begin{pmatrix} u_L & u_R \\ d_L & d_R \end{pmatrix}, \quad L = \begin{pmatrix} \nu_L & ? \\ e_L & e_R \end{pmatrix}$$

Action de \mathcal{A} sur \mathcal{H}_F

$$a = (\lambda, q, m) \in \mathcal{A}, \quad q = \begin{pmatrix} \alpha & \beta \\ -\bar{\beta} & \bar{\alpha} \end{pmatrix}$$

$$\begin{aligned} a u_R &= \lambda u_R & a u_L &= \alpha u_L - \bar{\beta} d_L \\ a d_R &= \bar{\lambda} d_R & a d_L &= \beta u_L + \bar{\alpha} d_L. \end{aligned}$$

$$a \bar{f} = \lambda \bar{f} \quad \text{si } f \text{ est un lepton}$$

$$a \bar{f} = m \bar{f} \quad \text{si } f \text{ est un quark}$$

$$\gamma(f_R) = f_R, \quad \gamma(f_L) = -f_L$$

Espace Fini

$$D_F = \begin{pmatrix} Y & 0 \\ 0 & \bar{Y} \end{pmatrix}$$

$$Y = Y_q \otimes 1_3 \oplus Y_\ell$$

$$Y_q = \begin{pmatrix} 0 & 0 & M_u & 0 \\ 0 & 0 & 0 & M_d \\ M_u^* & 0 & 0 & 0 \\ 0 & M_d^* & 0 & 0 \end{pmatrix}$$

$$Y_\ell = \begin{pmatrix} 0 & 0 & M_e \\ 0 & 0 & 0 \\ M_e^* & 0 & 0 \end{pmatrix}$$

Fluctuations internes \rightarrow Bosons

$$A = \sum a_i [D, a'_i] \quad a_i, a'_i \in \mathcal{A}$$

$$\sum a_i [\gamma_5 \otimes D_F, a'_i] \rightarrow \mathbf{Higgs}$$

$$\begin{pmatrix} 0 & X \\ X' & 0 \end{pmatrix}$$

$$X = \begin{pmatrix} M_u \varphi_1 & M_u \varphi_2 \\ -M_d \bar{\varphi}_2 & M_d \bar{\varphi}_1 \end{pmatrix}$$

$$X' = \begin{pmatrix} M_u^* \varphi'_1 & M_d^* \varphi'_2 \\ -M_u^* \bar{\varphi}'_2 & M_d^* \bar{\varphi}'_1 \end{pmatrix}$$

$\sum a_i [D_M \otimes 1, a'_i] \rightarrow$ **potentiels de jauge**

$$a_i = (\lambda_i, q_i, m_i), \quad a'_i = (\lambda'_i, q'_i, m'_i)$$

$U(1)$ potentiel de jauge $\Lambda = \sum \lambda_i d \lambda'_i$

$SU(2)$ potentiel de jauge $Q = \sum q_i d q'_i$

$U(3)$ potentiel de jauge $V = \sum m_i d m'_i$.

Hypercharges

$$D \mapsto \tilde{D} = D + A + JAJ^{-1} \quad A = A^*$$

$$\text{trace } V = \Lambda$$

$$V = V' + \frac{1}{3} \Lambda$$

V' est un potentiel de jauge $SU(3)$

$$\begin{pmatrix} \frac{4}{3} \Lambda + V' & 0 & 0 & 0 \\ 0 & -\frac{2}{3} \Lambda + V' & 0 & 0 \\ 0 & 0 & Q_{11} + \frac{1}{3} \Lambda + V' & Q_{12} \\ 0 & 0 & Q_{21} & Q_{22} + \frac{1}{3} \Lambda + V' \end{pmatrix}$$

$$\begin{pmatrix} -2\Lambda & 0 & 0 \\ 0 & Q_{11} - \Lambda & Q_{12} \\ 0 & Q_{21} & Q_{22} - \Lambda \end{pmatrix}$$

S. Coleman, *Renormalization and symmetry : a review for non-specialists*, in “Aspects of Symmetry” pp.99–112, Cambridge University Press, 1985.

A. Chamseddine and A. Connes, *The spectral action principle*, Commun. Math. Phys. 186 (1997) 731–750.

A. Chamseddine and A. Connes, *Scale Invariance in the spectral action*, Arxiv hep-th/0512169.

J. Collins, *Renormalization*, Cambridge Monographs in Math. Physics, Cambridge University Press, 1984.

A. Connes, M. Marcolli, *Renormalization and motivic Galois theory*, International Math. Research Notices, (2004), no. 76, 4073–4091.

A. Connes, M. Marcolli, *Anomalies, Dimensional Regularization and Noncommutative Geometry*, in preparation.

M. Veltman, *Diagrammatica : the path to Feynman diagrams*, Cambridge Univ. Press, 1994.